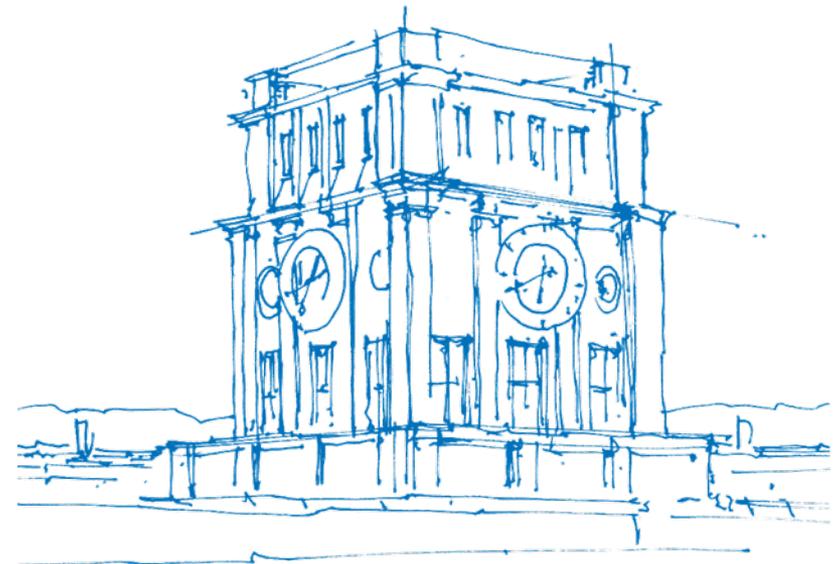


# Improving Time Stepping in Partitioned Multi-Physics

Benjamin R uth, Benjamin Uekermann

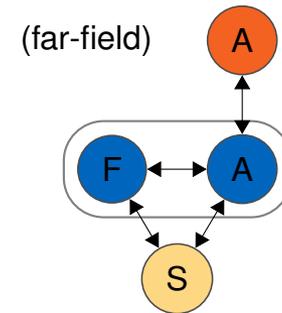
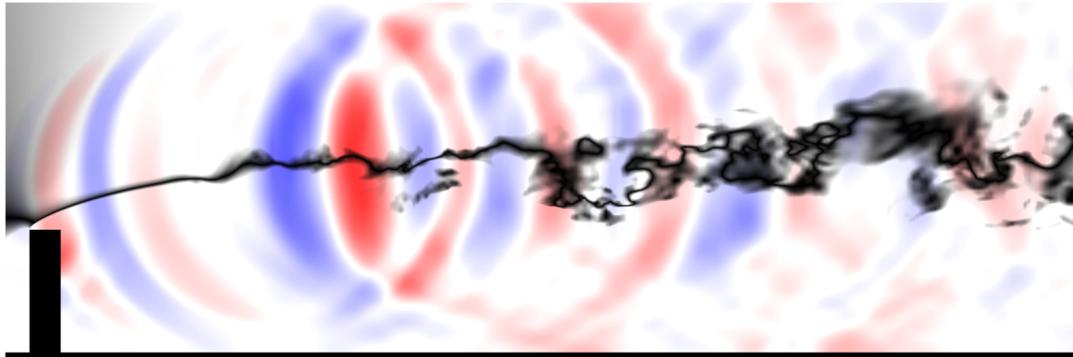
Technical University of Munich  
Department of Informatics  
Chair of Scientific Computing

89th GAMM Annual Meeting  
Technical University of Munich  
20. March 2018



*TUM Uhrenturm*

# Fluid-Structure-Acoustics

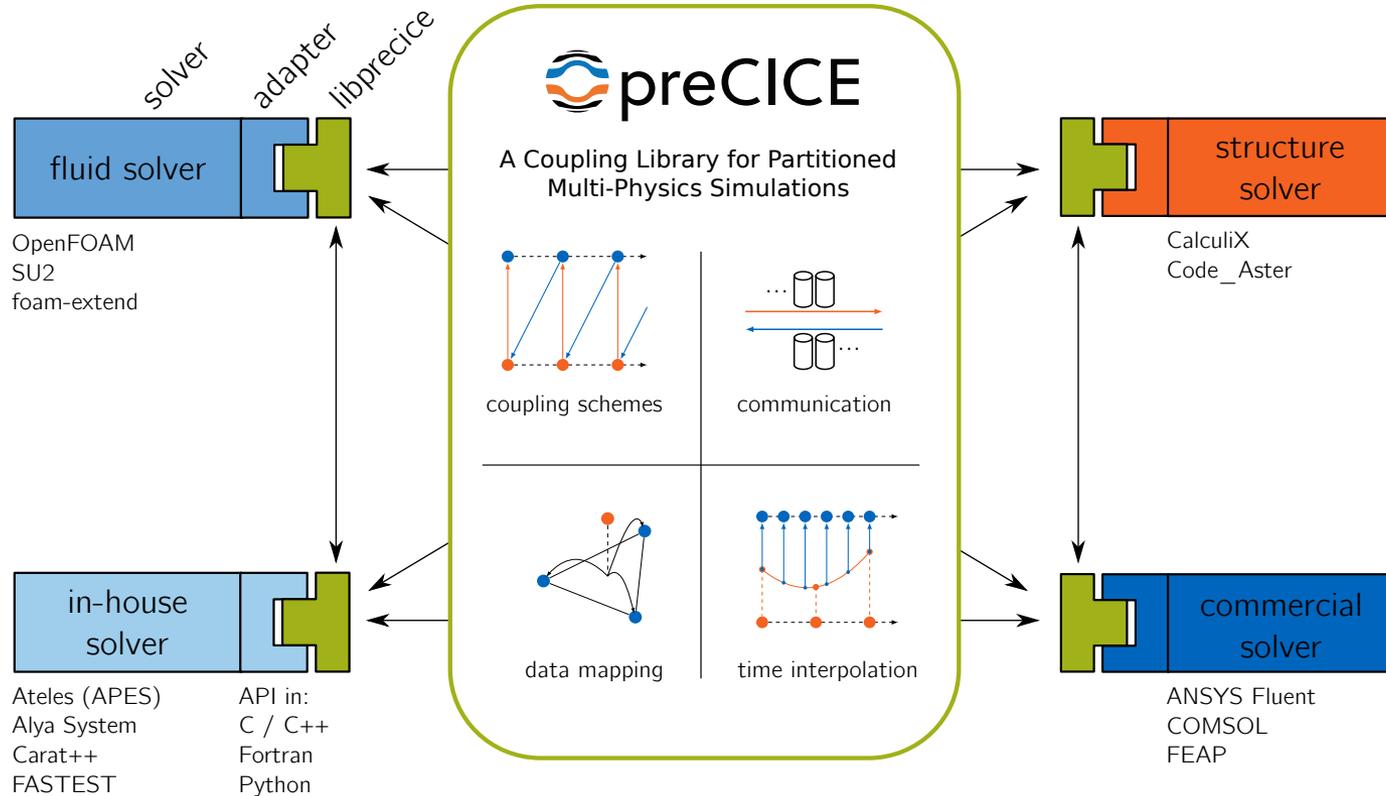


Fluid-Structure-Acoustics simulation and partitioned setup<sup>1</sup>.

physics	timescale	solver	scheme	order
(A)	small	Ateles	RK	2 or 4
(A)	small	FASTEST	EE	1
(F)	medium	FASTEST	CN	2
(S)	large	FEAP	$N\text{-}\beta$	1 or 2

<sup>1</sup>Reimann, T., et al. (2017). Aspects of FSI with aeroacoustics in turbulent flow. In 7th GACM Colloquium on Computational Mechanics.

# preCICE<sup>1</sup>



<sup>1</sup>Bungartz, H.-J., et al. (2016). preCICE – A fully parallel library for multi-physics surface coupling.

<https://doi.org/10.1016/j.compfluid.2016.04.003>

# preCICE at GAMM

## preCICE Coupling Library for Multi-Physics Simulation

Amin Totounferoush, University of Stuttgart S07.01 Coupled Problems  
(today in the morning)

## Quasi-Newton – A Universal Approach for Coupled Problems and Optimization

Miriam Mehl, University of Stuttgart S07.01 Coupled Problems  
(just now)

## Multi-physics simulations with OpenFOAM through preCICE

Gerasimos Chourdakis, Technical University of Munich S22.01 Scientific Computing  
(Thursday morning)

# Improving Time Stepping in Partitioned Multi-Physics

## Requirements

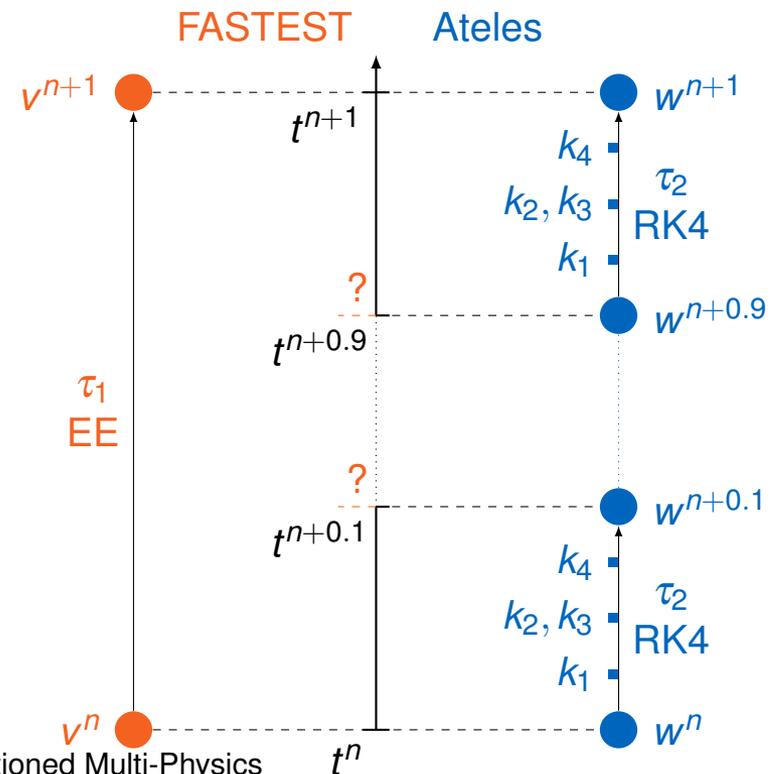
### Engineering:

- use different solvers
- use different discretization
- no degradation of solver performance

### Informatics:

- black-box approach
- parallel

### Multi-Scale Multi-Physics

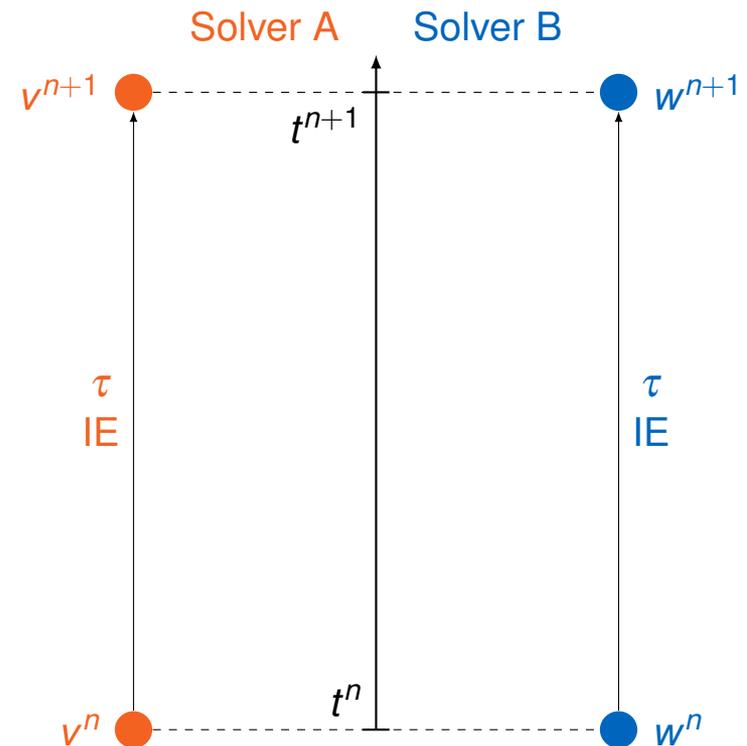


# Outline

## Partitioned Heat Transport Equation

- introduce the partitioned heat transport equation example
- introduce classical and advanced coupling schemes
- show deficits of classical explicit and implicit coupling schemes
- show advantages of waveform relaxation coupling scheme

## Simple setup



# Reference Solution: The Monolithic Setup

## Heat Transport equation

$$\frac{\partial u(x, t)}{\partial t} = \alpha \frac{\partial^2 u(x, t)}{\partial x^2}, x \in \Omega, t \in \mathbb{R}^+$$

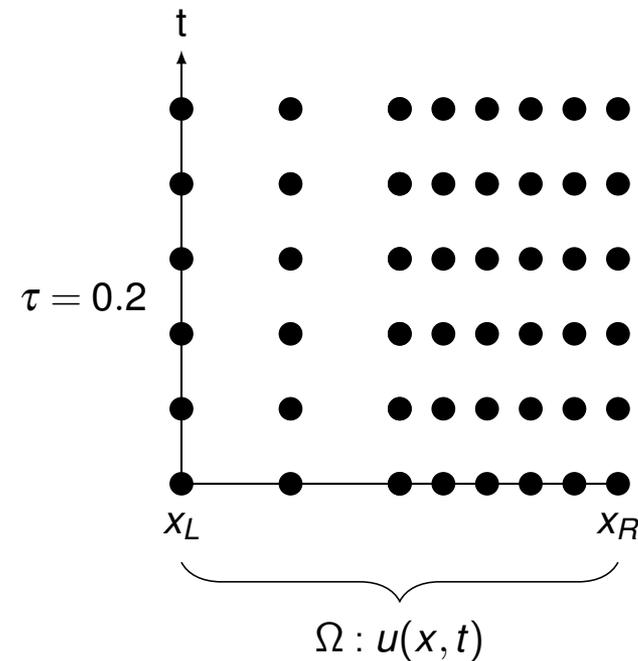
## Dirichlet boundary conditions

$$u(x = x_L, t) = u_L^D, u(x = x_R, t) = u_R^D$$

## Initial condition

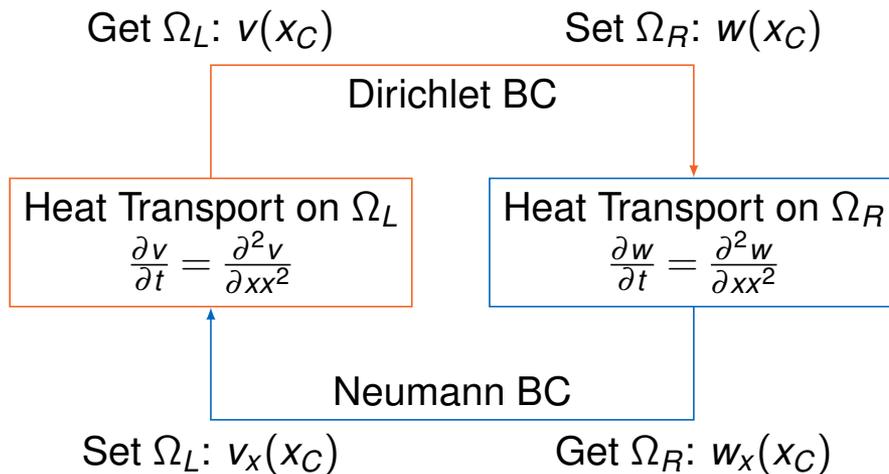
$$u(x, t = 0) = u_0(x)$$

## Discretization

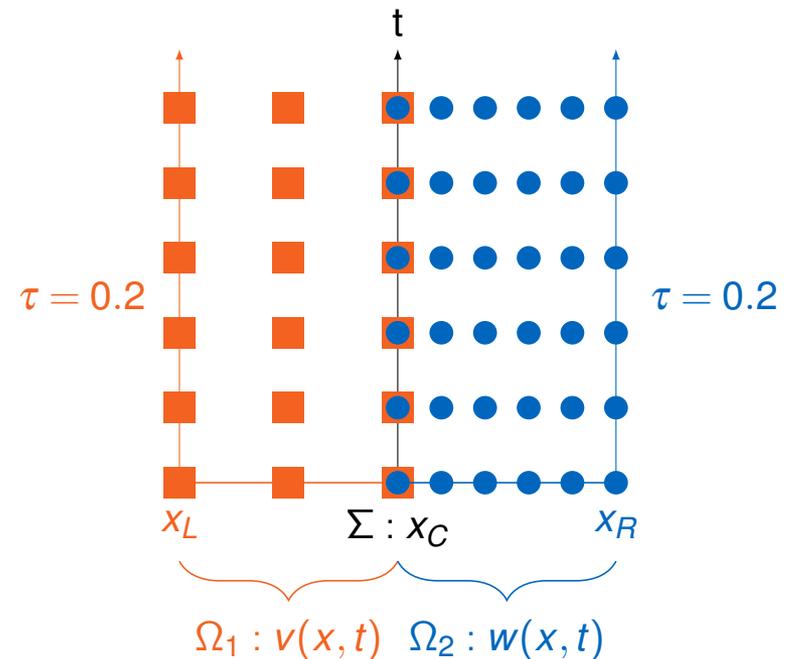


# The Partitioned Setup

## Dirichlet-Neumann coupling

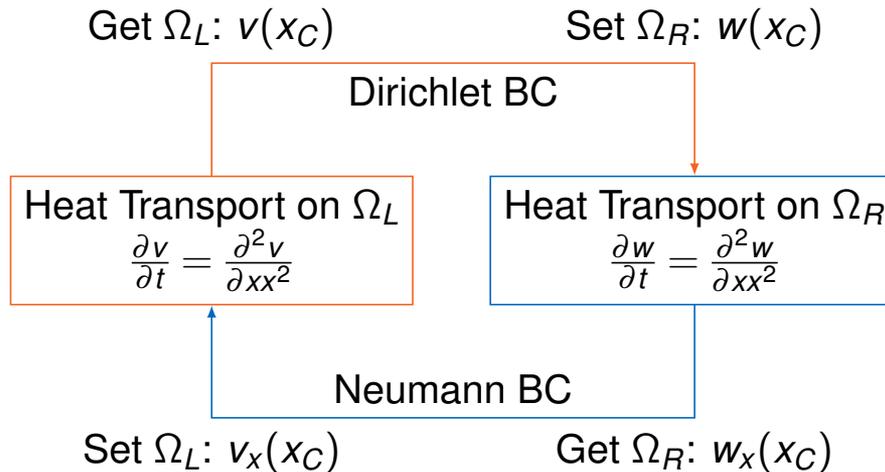


## Partitioning

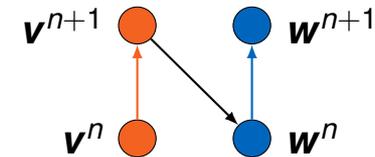


# Classical Coupling Schemes

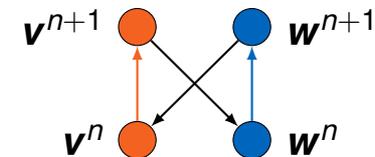
## Dirichlet-Neumann coupling



## Explicit coupling

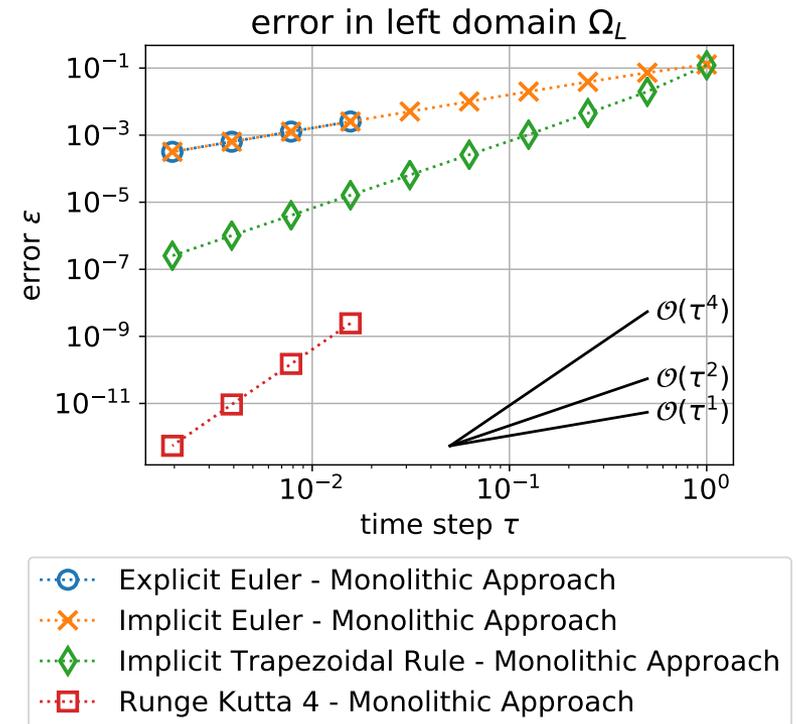


## Implicit coupling



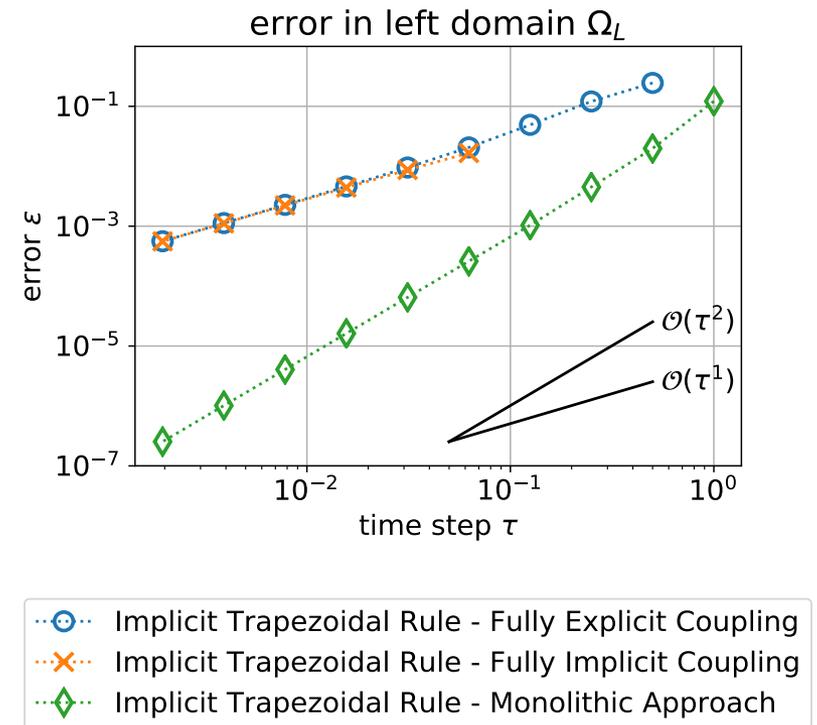
# Convergence order in time

- use constant spatial meshwidth  $h$
- refine temporal meshwidth  $\tau$
- compare to monolithic reference solution  $\mathbf{u}^n$  with fine  $\tau$



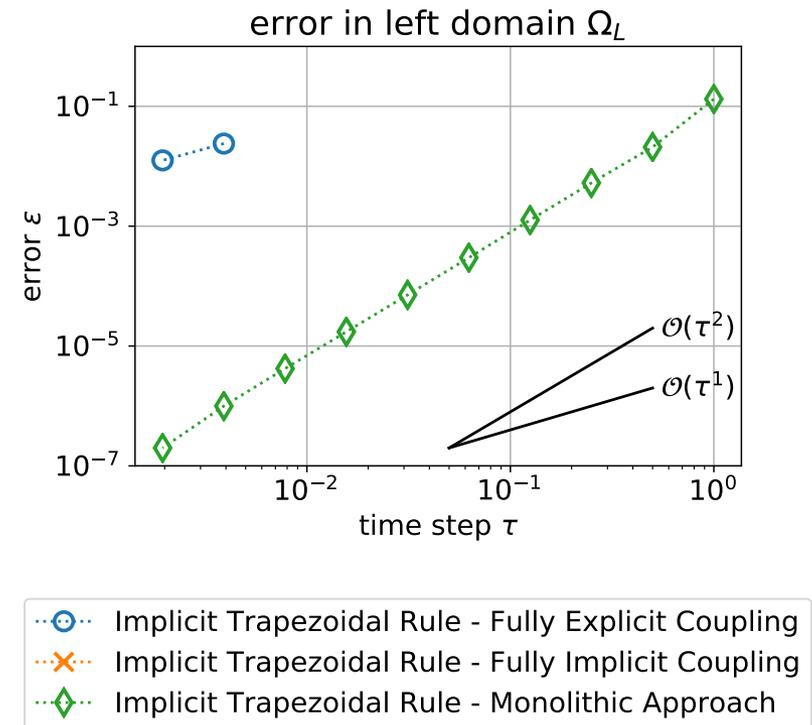
# Order Degradation: Trapezoidal rule

- order reduced to  $\mathcal{O}(\tau)$
- $h = 0.2$
- stability problems for Fully implicit coupling



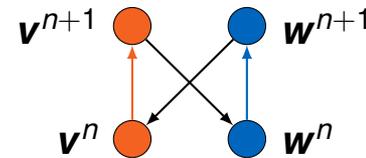
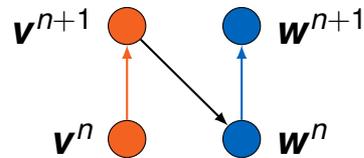
# Order Degradation: Trapezoidal rule

- order reduced to  $\mathcal{O}(\tau)$
- $h = 0.01$
- stability problems for Fully implicit coupling
- stability problems for Fully explicit coupling



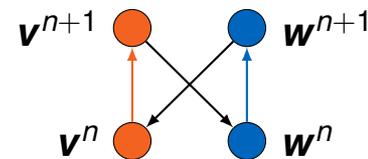
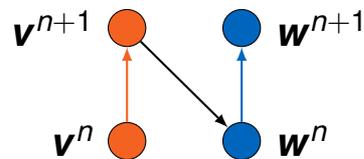
# Semi Implicit-Explicit Coupling

	update scheme	stability	order
fully explicit	$\mathbf{v}^{n+1} = \mathbf{v}^n + \frac{\tau}{2} [f_v(\mathbf{v}^n, t_n) + f_v(\mathbf{v}^{n+1}, t_{n+1})]$ $\mathbf{w}^{n+1} = \mathbf{w}^n + \frac{\tau}{2} [f_w(\mathbf{w}^n, t_n) + f_w(\mathbf{w}^{n+1}, t_{n+1})]$	depends on $\tau$	$\mathcal{O}(\tau)$
fully implicit	$\mathbf{v}^{n+1} = \mathbf{v}^n + \frac{\tau}{2} [f_v(\mathbf{v}^n, t_n) + f_v(\mathbf{v}^{n+1}, t_{n+1})]$ $\mathbf{w}^{n+1} = \mathbf{w}^n + \frac{\tau}{2} [f_w(\mathbf{w}^n, t_n) + f_w(\mathbf{w}^{n+1}, t_{n+1})]$	depends on $\tau$	$\mathcal{O}(\tau)$



# Semi Implicit-Explicit Coupling

	update scheme	stability	order
fully explicit	$\mathbf{v}^{n+1} = \mathbf{v}^n + \frac{\tau}{2} [f_v(\mathbf{v}^n, t_n, \mathbf{c}_n) + f_v(\mathbf{v}^{n+1}, t_{n+1}, \mathbf{c}_n)]$ $\mathbf{w}^{n+1} = \mathbf{w}^n + \frac{\tau}{2} [f_w(\mathbf{w}^n, t_n, \mathbf{c}_{n+1}) + f_w(\mathbf{w}^{n+1}, t_{n+1}, \mathbf{c}_{n+1})]$	depends on $\tau$	$\mathcal{O}(\tau)$
fully implicit	$\mathbf{v}^{n+1} = \mathbf{v}^n + \frac{\tau}{2} [f_v(\mathbf{v}^n, t_n, \mathbf{c}_{n+1}) + f_v(\mathbf{v}^{n+1}, t_{n+1}, \mathbf{c}_{n+1})]$ $\mathbf{w}^{n+1} = \mathbf{w}^n + \frac{\tau}{2} [f_w(\mathbf{w}^n, t_n, \mathbf{c}_{n+1}) + f_w(\mathbf{w}^{n+1}, t_{n+1}, \mathbf{c}_{n+1})]$	depends on $\tau$	$\mathcal{O}(\tau)$

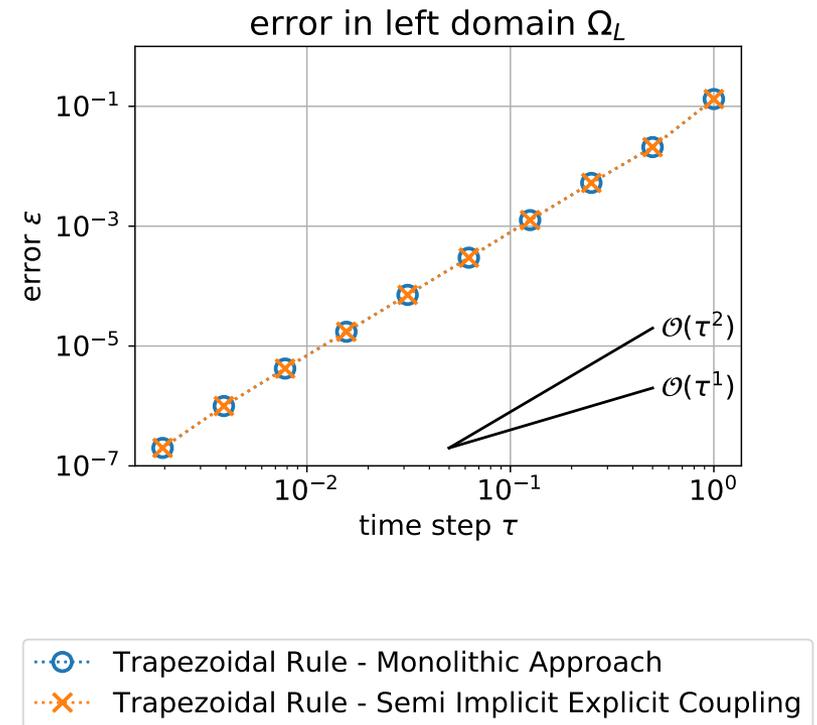


# Semi Implicit-Explicit Coupling

	update scheme	stability	order
fully explicit	$\mathbf{v}^{n+1} = \mathbf{v}^n + \frac{\tau}{2} [f_v(\mathbf{v}^n, t_n, \mathbf{c}_n) + f_v(\mathbf{v}^{n+1}, t_{n+1}, \mathbf{c}_n)]$ $\mathbf{w}^{n+1} = \mathbf{w}^n + \frac{\tau}{2} [f_w(\mathbf{w}^n, t_n, \mathbf{c}_{n+1}) + f_w(\mathbf{w}^{n+1}, t_{n+1}, \mathbf{c}_{n+1})]$	depends on $\tau$	$\mathcal{O}(\tau)$
fully implicit	$\mathbf{v}^{n+1} = \mathbf{v}^n + \frac{\tau}{2} [f_v(\mathbf{v}^n, t_n, \mathbf{c}_{n+1}) + f_v(\mathbf{v}^{n+1}, t_{n+1}, \mathbf{c}_{n+1})]$ $\mathbf{w}^{n+1} = \mathbf{w}^n + \frac{\tau}{2} [f_w(\mathbf{w}^n, t_n, \mathbf{c}_{n+1}) + f_w(\mathbf{w}^{n+1}, t_{n+1}, \mathbf{c}_{n+1})]$	depends on $\tau$	$\mathcal{O}(\tau)$
semi explicit-implicit	$\mathbf{v}^{n+1} = \mathbf{v}^n + \frac{\tau}{2} [f_v(\mathbf{v}^n, t_n, \mathbf{c}_n) + f_v(\mathbf{v}^{n+1}, t_{n+1}, \mathbf{c}_{n+1})]$ $\mathbf{w}^{n+1} = \mathbf{w}^n + \frac{\tau}{2} [f_w(\mathbf{w}^n, t_n, \mathbf{c}_n) + f_w(\mathbf{w}^{n+1}, t_{n+1}, \mathbf{c}_{n+1})]$	???	???

# Higher Order: Semi Implicit-Explicit Coupling

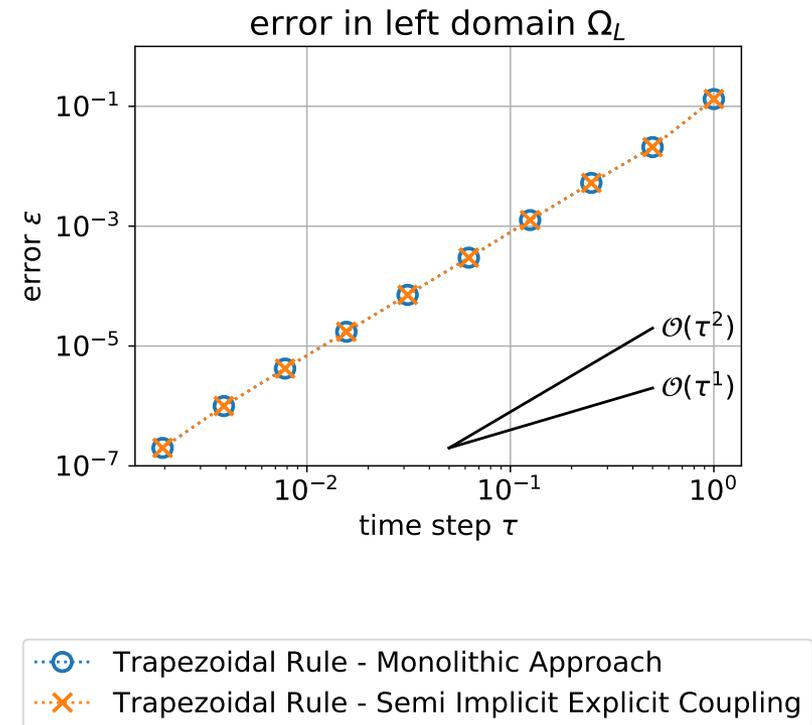
- order  $\mathcal{O}(\tau^2)$  maintained for semi explicit-implicit coupling
- no stability problems for semi explicit-implicit coupling



# Higher Order: Semi Implicit-Explicit Coupling

- order  $\mathcal{O}(\tau^2)$  maintained for semi explicit-implicit coupling
- no stability problems for semi explicit-implicit coupling

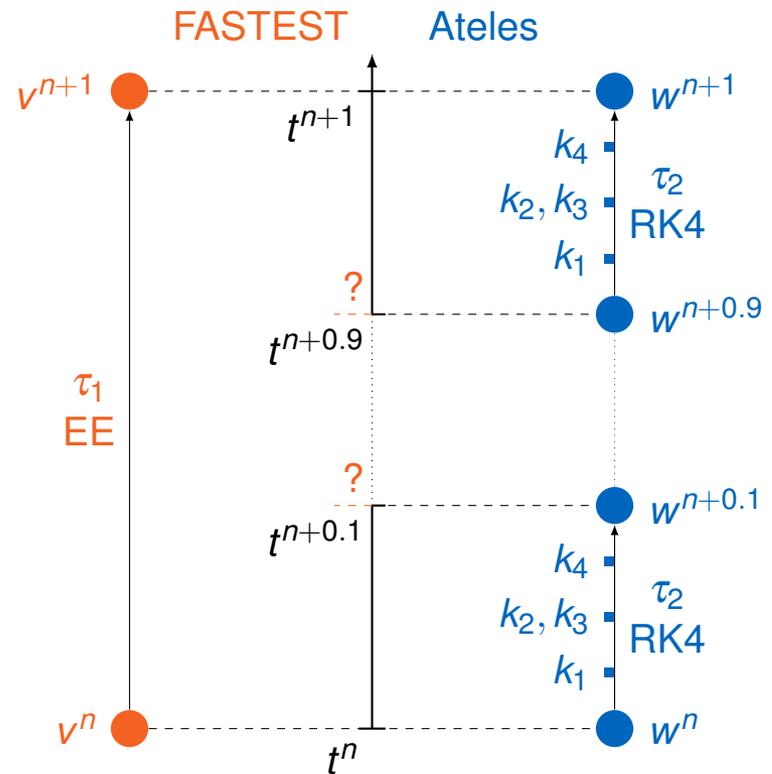
	stability	order
fully explicit	depends on $\tau$	$\mathcal{O}(\tau)$
fully implicit	depends on $\tau$	$\mathcal{O}(\tau)$
semi explicit-implicit	unconditionally	$\mathcal{O}(\tau^2)$



# Intermediate Summary

Semi-Explicit-Implicit	Goal
identical timesteps	subcycling
simple schemes	substepping
identical solvers	inhomogeneous setup
$\mathcal{O}(\tau^2)$	Higher order
taylorred schemes	general solution strategy

## Multi-Scale Multi-Physics



# What is Waveform Relaxation?

## Background Information

### Algorithm<sup>1</sup>

We want to solve the coupled problem

$$F_v(v, c) = 0, F_w(w, c) = 0.$$

with  $v, w, c$  known for  $t < t_n$  on the window  $T_n = [t_n, t_{n+1}]$ .

1. set  $k = 0$  and extrapolate  $c^0(t) = c_n$  for  $t \in T$
2. solve decoupled  $F_v, F_w$  using  $c^k$  to obtain  $v^{k+1}, w^{k+1}$  for  $t \in T$
3. use  $v^{k+1}, w^{k+1}$  to obtain  $c^{k+1}$
4. if not converged:
  - a. set  $k = k + 1$  and go to step 2,
  - b. otherwise proceed to next window  $T_{n+1}$

---

<sup>1</sup>Adapted from *Maciejewski, M., et al. (2017). Application of the Waveform Relaxation Technique to the Co-Simulation of Power Converter Controller and Electrical Circuit Models. <https://doi.org/10.1109/MMAR.2017.8046937>*

# Waveform Relaxation (WR) Coupling Scheme

## WR with our example

- Semi-Explicit-Implicit coupling equals WR with linear interpolation of

$$c^k(t) = \frac{c_n(t_{n+1} - t)}{\tau} + \frac{c_{n+1}^k(t - t_n)}{\tau}.$$

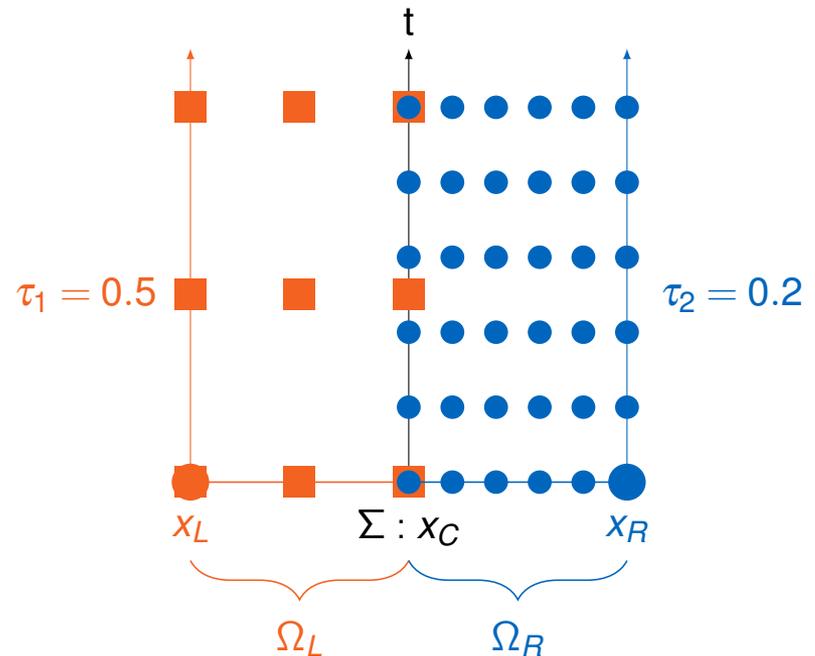
- Semi-Explicit-Implicit coupling:

$$\mathbf{v}^{n+1} = \mathbf{v}^n + \frac{\tau}{2} [f_v(c^k(t_n)) + f_v(c^k(t_{n+1}))]$$

$$\mathbf{w}^{n+1} = \mathbf{w}^n + \frac{\tau}{2} [f_w(c^k(t_n)) + f_w(c^k(t_{n+1}))]$$

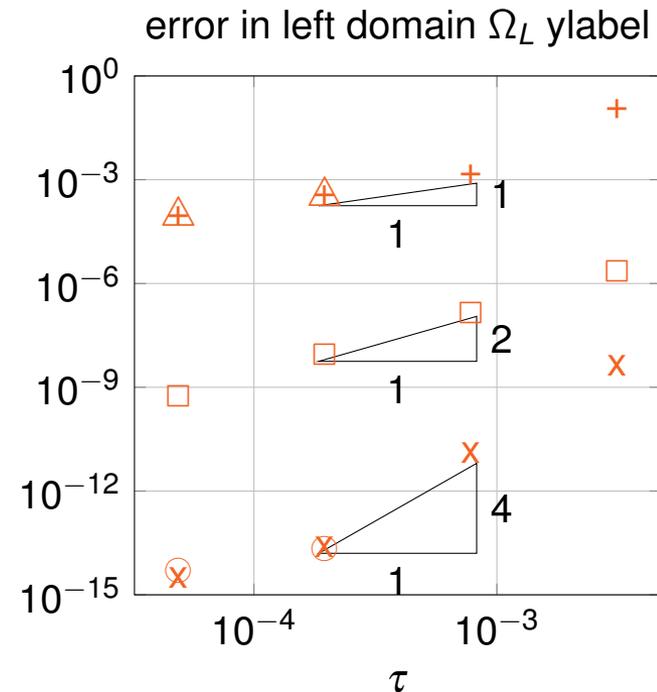
- Other interpolation methods are spline or dense output interpolation

## Multi-Scale Setup



# High Order and Subcycling

	scheme-solvers	time step	order	stable
○	M-RK4	$\tau_1 = \tau_2$	$\mathcal{O}(\tau^4)$	small $\tau_2$
+	Im-TR/TR	$\tau_1 = \tau_2$	$\mathcal{O}(\tau)$	small $\tau_2$
△	Im-RK4/RK4	$\tau_1 = \tau_2$	$\mathcal{O}(\tau)$	small $\tau_2$
x	WR-RK4/RK4	$\tau_1 > \tau_2$	$\mathcal{O}(\tau^4)$	small $\tau_2$
□	WR-RK4/TR	$\tau_1 = \tau_2$	$\mathcal{O}(\tau^2)$	$\forall \tau_2$

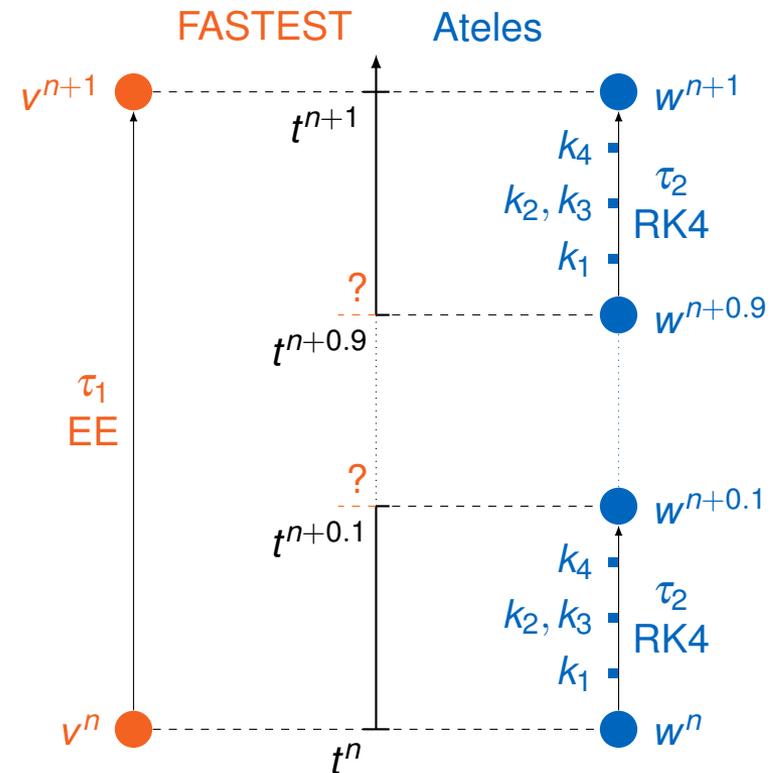


# Conclusion

## Partitioned Heat Transport

- ✓ introduce the partitioned heat transport equation example
- ✓ introduce classical and advanced coupling schemes
- ✓ deficits of classical explicit and implicit coupling schemes
  - order and stability degradation
- ✓ advantages of waveform relaxation coupling scheme
  - order and stability maintained
  - inhomogeneous setup
  - subcycling

## Multi-Scale



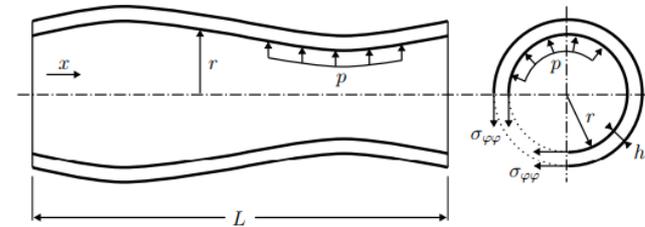
# Outlook

## Implementation

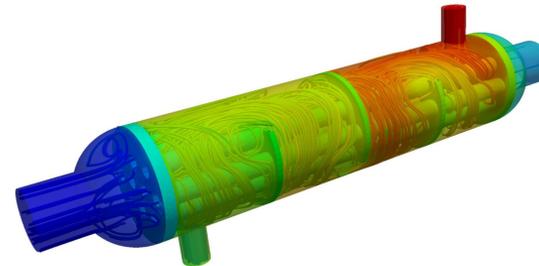
- Interpolation methods?
- Convergence of acceleration schemes
- Parallel performance

## Further Tests

### 1D Tube<sup>1</sup>:



### preCICE examples<sup>2</sup>:



<sup>1</sup>figure from Degroote, J., et al. (2008). Stability of a coupling technique for partitioned solvers in FSI applications. <https://doi.org/10.1016/j.compstruc.2008.05.005>

<sup>2</sup>figure from Cheung Yau, L. (2016). Conjugate Heat Transfer with the Multiphysics Coupling Library preCICE. TUM.