

A novel derivation for modal derivatives based on Volterra series representation and its use in nonlinear model order reduction

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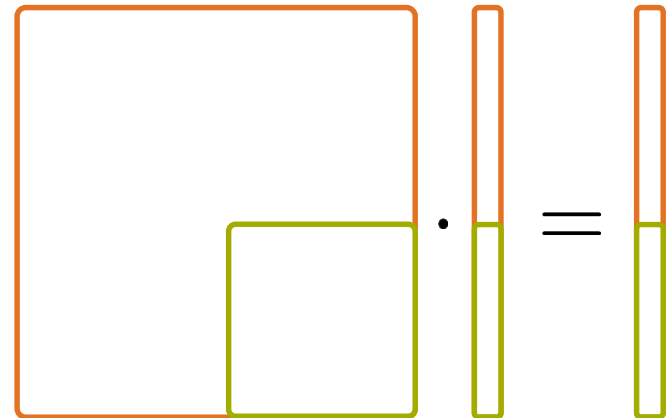
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COMPDYN-UNCECOMP 2019

MS 27: Advances in MOR techniques for CSD

Crete, June 26th 2019

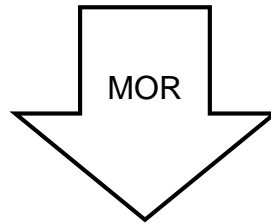


Projective Model Order Reduction

Nonlinear second-order full order model (FOM)

$$M\ddot{q}(t) + D\dot{q}(t) + f(q(t)) = BF(t) \quad q(0) = q_0, \dot{q}(0) = \dot{q}_0,$$
$$y(t) = Cq(t)$$

Linear Galerkin projection



$$q(t) \approx Vq_r(t), \quad V \in \mathbb{R}^{n \times r} \quad r \ll n$$

Reduced order model (ROM)

$$M_r\ddot{q}_r(t) + D_r\dot{q}_r(t) + V^T f(Vq_r(t)) = B_r F(t) \quad \{q_r(0), \dot{q}_r(0)\} = (V^T V)^{-1} V^T \{q_0, \dot{q}_0\},$$
$$y_r(t) = C_r q_r(t)$$

with

$$\{M_r, D_r\} = V^T \{M, D\} V, \quad f_r(q) = V^T f(Vq_r) \quad \text{Hyper-reduction!}$$

$$B_r = V^T B,$$

$$C_r = C V,$$

In this talk: Dimensional reduction
How to choose V ?

Simulation-based approaches (e.g. POD)

Take snapshots of the simulated trajectory for typical (training) input force and perform SVD

$$\underset{(n, n_s)}{Q} = [\mathbf{q}(t_1), \mathbf{q}(t_2), \dots, \mathbf{q}(t_{n_s})]$$

Reduction basis: $V = M_r \in \mathbb{R}^{n \times r}$

$$Q \stackrel{\text{SVD}}{=} \underset{(n, n)}{M} \underset{(n, n_s)}{\Sigma} \underset{(n_s, n_s)}{N^T} \approx \underset{(n, r)}{M_r} \underset{(r, n_s)}{\Sigma_r} \underset{(n_s, n_s)}{N_r^T}$$

$$\mathbf{q}(t) \approx V \mathbf{q}_r(t) = \sum_{i=1}^r \mathbf{v}_i q_{r,i}(t)$$

Simulation-free / System-theoretic methods

- **Basis augmentation:** Enrichment of a linear basis with nonlinear information

$$V_{\text{aug}} = [V^{(1)}, V^{(2)}]$$

$$\mathbf{q}(t) \approx V_{\text{aug}} \mathbf{q}_{r, \text{aug}}(t)$$

- + : Easy projection
- : Higher reduced order

- **Nonlinear projection (e.g. Quadratic Manifold)**

$$V^{(1)} \in \mathbb{R}^{n \times r}$$

$$V^{(2)} \in \mathbb{R}^{n \times r^2}$$

$$\mathbf{q}(t) \approx V^{(1)} \mathbf{q}_r(t) + V^{(2)} (\mathbf{q}_r(t) \otimes \mathbf{q}_r(t))$$

- + : Smaller reduced order
- : Difficult projection

Reduced coordinates: $\mathbf{q}_r(t) = [q_{r,1}(t), \dots, q_{r,r}(t)]^T = [\eta_1(t), \dots, \eta_r(t)]^T$

Original derivation for modal derivatives

Idea: Compose reduction basis with both *vibration modes* and *modal derivatives*

$$\mathbf{V}_{\text{aug}} = [\mathbf{\Phi}_r, \mathbf{\Theta}_{r^2}] \quad \mathbf{\Phi}_r = [\phi_{1,\text{eq}}, \dots, \phi_{r,\text{eq}}] \in \mathbb{R}^{n \times r} \quad \mathbf{\Theta}_{r^2} = [\theta_{11}, \dots, \theta_{rr}] \in \mathbb{R}^{n \times r^2}$$

1.) Vibration modes of the linearized model:

Linearization point: \mathbf{q}_{eq}

$$(\mathbf{K}_{\text{eq}} - \omega_{i,\text{eq}}^2 \mathbf{M}) \phi_{i,\text{eq}} = \mathbf{0} \quad i = 1, \dots, r$$

$$\mathbf{K}_{\text{eq}} = \mathbf{K}(\mathbf{q}_{\text{eq}}) = \left. \frac{\partial \mathbf{f}(\mathbf{q}(t))}{\partial \mathbf{q}(t)} \right|_{\mathbf{q}_{\text{eq}}}$$

Normalization condition: $\phi_{i,\text{eq}}^\top \mathbf{M} \phi_{i,\text{eq}} = 1$

2.) Perturbation of eigenmodes:

$$\frac{\partial}{\partial \eta_j(t)} \left(\mathbf{K}(\mathbf{q}_{\text{eq}}) - \omega_i^2(\mathbf{q}_{\text{eq}}) \mathbf{M} \right) \phi_i(\mathbf{q}_{\text{eq}}) = \mathbf{0}$$



$$\left(\frac{\partial \mathbf{K}_{\text{eq}}}{\partial \eta_j(t)} - \frac{\partial \omega_{i,\text{eq}}^2}{\partial \eta_j(t)} \mathbf{M} \right) \phi_{i,\text{eq}} + \left(\mathbf{K}_{\text{eq}} - \omega_{i,\text{eq}}^2 \mathbf{M} \right) \frac{\partial \phi_{i,\text{eq}}}{\partial \eta_j(t)} = \mathbf{0} \quad \begin{array}{l} i = 1, \dots, r \\ j = 1, \dots, r \end{array}$$

Modal derivatives (MDs): $\theta_{ij} = \frac{\partial \phi_{i,\text{eq}}}{\partial \eta_j(t)}$

Calculation formula

$$\theta_{ij} = \frac{\partial \phi_{i,\text{eq}}}{\partial \eta_j(t)}$$

$$\left(\mathbf{K}_{\text{eq}} - \omega_{i,\text{eq}}^2 \mathbf{M} \right) \theta_{ij} = \left(\frac{\partial \omega_{i,\text{eq}}^2}{\partial \eta_j(t)} \mathbf{M} - \frac{\partial \mathbf{K}_{\text{eq}}}{\partial \eta_j(t)} \right) \phi_{i,\text{eq}}$$

$$\begin{aligned} i &= 1, \dots, r \\ j &= 1, \dots, r \end{aligned}$$

Singular linear system of equations

$$\mathbf{A} \mathbf{x} = \mathbf{b}$$

Right-hand side (rhs)

- Derivative of eigenfrequencies: $\phi_{i,\text{eq}}^\top \mathbf{M} \phi_{i,\text{eq}} = 1 \implies \frac{\partial \omega_{i,\text{eq}}^2}{\partial \eta_j(t)} = \phi_{i,\text{eq}}^\top \frac{\partial \mathbf{K}_{\text{eq}}}{\partial \eta_j(t)} \phi_{i,\text{eq}}$

- Finite difference scheme for tangential stiffness matrix:

- Forward difference: $\left. \frac{\partial \mathbf{K}(\mathbf{q})}{\partial \eta_j(t)} \right|_{\mathbf{q}_{\text{eq}}} = \frac{\mathbf{K}(\mathbf{q}_{\text{eq}} + \phi_{j,\text{eq}} \cdot h) - \mathbf{K}(\mathbf{q}_{\text{eq}})}{h}$

- Backward difference: $\left. \frac{\partial \mathbf{K}(\mathbf{q})}{\partial \eta_j(t)} \right|_{\mathbf{q}_{\text{eq}}} = \frac{\mathbf{K}(\mathbf{q}_{\text{eq}}) - \mathbf{K}(\mathbf{q}_{\text{eq}} - \phi_{j,\text{eq}} \cdot h)}{h}$

- Central difference: $\left. \frac{\partial \mathbf{K}(\mathbf{q})}{\partial \eta_j(t)} \right|_{\mathbf{q}_{\text{eq}}} = \frac{\mathbf{K}(\mathbf{q}_{\text{eq}} + \phi_{j,\text{eq}} \cdot h) - \mathbf{K}(\mathbf{q}_{\text{eq}} - \phi_{j,\text{eq}} \cdot h)}{2h}$

Handling the singular left-hand side (lhs)

By imposing an additional condition/constraint: $\frac{\partial}{\partial \eta_j(t)} \left(\phi_{i,\text{eq}}^\top M \phi_{i,\text{eq}} \right) = 0 \implies \boxed{\phi_{i,\text{eq}}^\top M \frac{\partial \phi_{i,\text{eq}}}{\partial \eta_j(t)} = 0}$

- Nelson's method: [\[Nelson '76\]](#)

- Direct method:
$$\begin{bmatrix} (\mathbf{K}_{\text{eq}} - \omega_{i,\text{eq}}^2 \mathbf{M}) & -\mathbf{M} \phi_{i,\text{eq}} \\ -(\mathbf{M} \phi_{i,\text{eq}})^\top & 0 \end{bmatrix} \begin{bmatrix} \theta_{ij} \\ \frac{\partial \omega_{i,\text{eq}}^2}{\partial \eta_j(t)} \end{bmatrix} = \begin{bmatrix} -\frac{\partial \mathbf{K}_{\text{eq}}}{\partial \eta_j(t)} \phi_{i,\text{eq}} \\ 0 \end{bmatrix} \quad \begin{array}{l} i = 1, \dots, r \\ j = 1, \dots, r \end{array}$$

Excluding mass consideration

Mass terms are usually neglected, leading to the so-called *static modal derivatives* (SMDs)

$$\left(\cancel{\mathbf{K}_{\text{eq}}} - \cancel{\omega_{i,\text{eq}}^2} \cancel{\mathbf{M}} \right) \frac{\partial \phi_{i,\text{eq}}}{\partial \eta_j(t)} = \left(\cancel{\frac{\partial \omega_{i,\text{eq}}^2}{\partial \eta_j(t)}} \cancel{\mathbf{M}} - \frac{\partial \mathbf{K}_{\text{eq}}}{\partial \eta_j(t)} \right) \phi_{i,\text{eq}}$$

$$\implies \boxed{\mathbf{K}_{\text{eq}} \frac{\partial \phi_{i,\text{eq}}}{\partial \eta_j(t)} \Big|_s = -\frac{\partial \mathbf{K}_{\text{eq}}}{\partial \eta_j(t)} \phi_{i,\text{eq}}} \quad \begin{array}{l} i = 1, \dots, r \\ j = 1, \dots, r \end{array}$$

Properties:

- Only one factorization of \mathbf{K}_{eq} is needed.

- Static modal derivatives are *symmetric*: $\theta_{s,ij} = \frac{\partial \phi_{i,\text{eq}}}{\partial \eta_j(t)} \Big|_s = \frac{\partial \phi_{j,\text{eq}}}{\partial \eta_i(t)} \Big|_s = \theta_{s,ji}$.

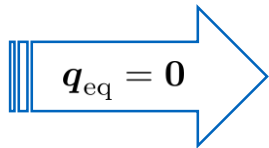
Polynomial system representation

Nonlinear second-order model

$$\begin{aligned} M\ddot{\mathbf{q}}(t) + D\dot{\mathbf{q}}(t) + \mathbf{f}(\mathbf{q}(t)) &= \mathbf{BF}(t) & \mathbf{q}(0) = \mathbf{q}_0, \dot{\mathbf{q}}(0) = \dot{\mathbf{q}}_0, \\ \mathbf{y}(t) &= \mathbf{C}\mathbf{q}(t) \end{aligned}$$

Taylor series expansion

$$\mathbf{f}(\mathbf{q}) = \mathbf{f}(\mathbf{q}_{\text{eq}}) + \frac{\partial \mathbf{f}(\mathbf{q}_{\text{eq}})}{\partial \mathbf{q}} (\mathbf{q} - \mathbf{q}_{\text{eq}}) + \frac{1}{2!} \frac{\partial^2 \mathbf{f}(\mathbf{q}_{\text{eq}})}{\partial \mathbf{q}^2} (\mathbf{q} - \mathbf{q}_{\text{eq}})^{(2)} + \frac{1}{3!} \frac{\partial^3 \mathbf{f}(\mathbf{q}_{\text{eq}})}{\partial \mathbf{q}^3} (\mathbf{q} - \mathbf{q}_{\text{eq}})^{(3)} + \dots$$



$$\mathbf{f}(\mathbf{q}) = \mathbf{K}^{(1)} \mathbf{q} + \mathbf{K}^{(2)} (\mathbf{q} \otimes \mathbf{q}) + \mathbf{K}^{(3)} (\mathbf{q} \otimes \mathbf{q} \otimes \mathbf{q}) + \dots$$

Polynomial (cubic) nonlinear second-order model

$$M\ddot{\mathbf{q}}(t) + D\dot{\mathbf{q}}(t) + \mathbf{K}^{(1)} \mathbf{q}(t) + \mathbf{K}^{(2)} (\mathbf{q}(t) \otimes \mathbf{q}(t)) + \mathbf{K}^{(3)} (\mathbf{q}(t) \otimes \mathbf{q}(t) \otimes \mathbf{q}(t)) = \mathbf{BF}(t)$$

Symmetric tensors $\mathbf{K}^{(1)} \in \mathbb{R}^{n \times n}$, $\mathcal{K}^{(2)} \in \mathbb{R}^{n \times n \times n}$ and $\mathcal{K}^{(3)} \in \mathbb{R}^{n \times n \times n \times n}$

$$\mathbf{K}_{ab}^{(1)} = \frac{\partial f_a}{\partial q_b} = \frac{\partial^2 \mathcal{V}}{\partial q_a \partial q_b}, \quad \mathcal{K}_{abc}^{(2)} = \frac{1}{2} \frac{\partial^2 f_a}{\partial q_b \partial q_c} = \frac{1}{2} \frac{\partial^3 \mathcal{V}}{\partial q_a \partial q_b \partial q_c}, \quad \mathcal{K}_{abcd}^{(3)} = \frac{1}{6} \frac{\partial^3 f_a}{\partial q_b \partial q_c \partial q_d} = \frac{1}{6} \frac{\partial^4 \mathcal{V}}{\partial q_a \partial q_b \partial q_c \partial q_d}$$

Volterra theory and variational equations

For an input of the form $\alpha F(t)$, it is assumed that the response is

$$\begin{aligned} \mathbf{q}(t) &= \alpha \mathbf{q}_1(t) + \alpha^2 \mathbf{q}_2(t) + \dots, \\ \ddot{\mathbf{q}}(t) &= \alpha \ddot{\mathbf{q}}_1(t) + \alpha^2 \ddot{\mathbf{q}}_2(t) + \dots. \end{aligned}$$

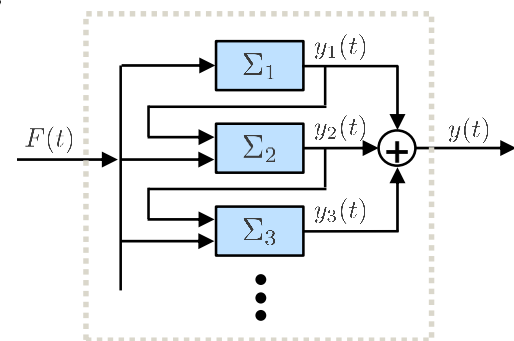
aka. Poincaré expansion
[Nayfeh '08], [Rugh '81]

Inserting the assumed input and the assumed response into the polynomial system, yields

$$\begin{aligned} M(\alpha \ddot{\mathbf{q}}_1(t) + \alpha^2 \ddot{\mathbf{q}}_2(t) + \dots) &+ \mathbf{K}^{(1)}(\alpha \mathbf{q}_1(t) + \alpha^2 \mathbf{q}_2(t) + \dots) \\ &+ \mathbf{K}^{(2)}(\mathbf{q}(t) \otimes \mathbf{q}(t)) + \mathbf{K}^{(3)}(\mathbf{q}(t) \otimes \mathbf{q}(t) \otimes \mathbf{q}(t)) = \mathbf{B} \alpha F(t). \end{aligned}$$

Equating coefficients of α^k , yields the *variational equations* ($D = 0$)

$$\begin{aligned} \alpha : \quad M \ddot{\mathbf{q}}_1(t) + \mathbf{K}^{(1)} \mathbf{q}_1(t) &= \mathbf{B} F(t), & \mathbf{q}_1(0) &= \mathbf{q}_0, \\ \alpha^2 : \quad M \ddot{\mathbf{q}}_2(t) + \mathbf{K}^{(1)} \mathbf{q}_2(t) &= -\mathbf{K}^{(2)}(\mathbf{q}_1(t) \otimes \mathbf{q}_1(t)), & \mathbf{q}_2(0) &= \mathbf{0}, \\ \alpha^3 : \quad M \ddot{\mathbf{q}}_3(t) + \mathbf{K}^{(1)} \mathbf{q}_3(t) &= -\mathbf{K}^{(2)}(\mathbf{q}_1(t) \otimes \mathbf{q}_2(t) + \mathbf{q}_2(t) \otimes \mathbf{q}_1(t)) - \mathbf{K}^{(3)}(\mathbf{q}_1(t) \otimes \mathbf{q}_1(t) \otimes \mathbf{q}_1(t)), & \mathbf{q}_3(0) &= \mathbf{0}, \end{aligned}$$



series of cascaded subsystems

Modes (First subsystem)

Subsystem state-equation: $M\ddot{\mathbf{q}}_1(t) + \mathbf{K}^{(1)}\mathbf{q}_1(t) = \mathbf{B}F(t)$

Ansatz for the homogeneous solution: $\mathbf{q}_1(t) = \sum_{i=1}^n c_i \boldsymbol{\phi}_i \cos(\omega_i t)$

Inserting ansatz (with $\ddot{\mathbf{q}}_1(t)$) in state-equation yields: $(\mathbf{K}^{(1)} - \omega_i^2 \mathbf{M}) \boldsymbol{\phi}_i \sum_{i=1}^n c_i \cos(\omega_i t) = \mathbf{0}$,

Modal derivatives (Second subsystem)

Subsystem state-equation: $M\ddot{\mathbf{q}}_2(t) + \mathbf{K}^{(1)}\mathbf{q}_2(t) = -\mathbf{K}^{(2)}(\mathbf{q}_1(t) \otimes \mathbf{q}_1(t))$

Ansatz for the particular solution: $\mathbf{q}_2(t) = \sum_{i=1}^n \sum_{j=1}^n \frac{1}{2} c_i c_j \left(\tilde{\boldsymbol{\theta}}_{ij} \cos((\omega_i + \omega_j)t) + \tilde{\tilde{\boldsymbol{\theta}}}_{ij} \cos((\omega_i - \omega_j)t) \right)$
(method of undetermined coeff.)

Inserting ansatz (with $\ddot{\mathbf{q}}_2(t)$ and $\mathbf{q}_1(t) \otimes \mathbf{q}_1(t)$) in state-equation yields exemplarily:

$$\mathbf{0} = \frac{1}{2} c_1^2 \cos(2\omega_1 t) \underbrace{\left(\left(-(2\omega_1)^2 \mathbf{M} + \mathbf{K}^{(1)} \right) \tilde{\boldsymbol{\theta}}_{11} + \mathbf{K}^{(2)}(\boldsymbol{\phi}_1 \otimes \boldsymbol{\phi}_1) \right)}_{=0} + \frac{1}{2} c_1^2 \underbrace{\left(\mathbf{K}^{(1)} \tilde{\tilde{\boldsymbol{\theta}}}_{11} + \mathbf{K}^{(2)}(\boldsymbol{\phi}_1 \otimes \boldsymbol{\phi}_1) \right)}_{=0} + \dots$$

$$+ \frac{1}{2} c_1 c_2 \cos((\omega_1 + \omega_2)t) \underbrace{\left(\left(-(\omega_1 + \omega_2)^2 \mathbf{M} + \mathbf{K}^{(1)} \right) (\tilde{\boldsymbol{\theta}}_{12} + \tilde{\boldsymbol{\theta}}_{21}) + \mathbf{K}^{(2)}(\boldsymbol{\phi}_1 \otimes \boldsymbol{\phi}_2 + \boldsymbol{\phi}_2 \otimes \boldsymbol{\phi}_1) \right)}_{=0} + \dots$$

New modal derivatives

It follows from all brackets:

$$\begin{aligned} \left(\mathbf{K}^{(1)} - (\omega_i + \omega_j)^2 \mathbf{M} \right) \tilde{\boldsymbol{\theta}}_{ij} &= -\mathbf{K}^{(2)} (\boldsymbol{\phi}_i \otimes \boldsymbol{\phi}_j), & i, j = 1, \dots, r, \\ \left(\mathbf{K}^{(1)} - (\omega_i - \omega_j)^2 \mathbf{M} \right) \tilde{\tilde{\boldsymbol{\theta}}}_{ij} &= -\mathbf{K}^{(2)} (\boldsymbol{\phi}_i \otimes \boldsymbol{\phi}_j), & i, j = 1, \dots, r. \end{aligned}$$

Equivalent description for the right-hand side

$$\frac{\partial K_{ab}^{(1)}}{\partial \eta_j(t)} (\boldsymbol{\phi}_i)_b = \frac{\partial K_{ab}^{(1)}}{\partial q_c} \frac{\partial q_c}{\partial \eta_j(t)} (\boldsymbol{\phi}_i)_b := \frac{\partial^2 f_a}{\partial q_b \partial q_c} (\boldsymbol{\phi}_j)_c (\boldsymbol{\phi}_i)_b = 2 \mathcal{K}_{abc}^{(2)} (\boldsymbol{\phi}_j)_c (\boldsymbol{\phi}_i)_b.$$

$$\left. \frac{\partial \mathbf{K}^{(1)}(\mathbf{q})}{\partial \eta_j(t)} \right|_{\mathbf{q}_{\text{eq}}} \boldsymbol{\phi}_i := 2 \mathbf{K}^{(2)} (\boldsymbol{\phi}_i \otimes \boldsymbol{\phi}_j) = \mathbf{K}^{(2)} (\boldsymbol{\phi}_i \otimes \boldsymbol{\phi}_j + \boldsymbol{\phi}_j \otimes \boldsymbol{\phi}_i).$$

Thus, the *new modal derivatives* are given by:

$$\begin{aligned} \left(\mathbf{K}^{(1)} - (\omega_i + \omega_j)^2 \mathbf{M} \right) \tilde{\boldsymbol{\theta}}_{ij} &= -\frac{1}{2} \left. \frac{\partial \mathbf{K}^{(1)}(\mathbf{q})}{\partial \eta_j(t)} \right|_{\mathbf{q}_{\text{eq}}} \boldsymbol{\phi}_i, & i, j = 1, \dots, r, \\ \left(\mathbf{K}^{(1)} - (\omega_i - \omega_j)^2 \mathbf{M} \right) \tilde{\tilde{\boldsymbol{\theta}}}_{ij} &= -\frac{1}{2} \left. \frac{\partial \mathbf{K}^{(1)}(\mathbf{q})}{\partial \eta_j(t)} \right|_{\mathbf{q}_{\text{eq}}} \boldsymbol{\phi}_i, & i, j = 1, \dots, r. \end{aligned}$$

Conventional modal derivatives

$$\left(\mathbf{K}_{\text{eq}} - \omega_{i,\text{eq}}^2 \mathbf{M} \right) \boldsymbol{\theta}_{ij} = \left(\frac{\partial \omega_{i,\text{eq}}^2}{\partial \eta_j(t)} \mathbf{M} - \frac{\partial \mathbf{K}_{\text{eq}}}{\partial \eta_j(t)} \right) \boldsymbol{\phi}_{i,\text{eq}}$$

$$\mathbf{K}_{\text{eq}} \boldsymbol{\theta}_{s,ij} = - \frac{\partial \mathbf{K}_{\text{eq}}}{\partial \eta_j(t)} \boldsymbol{\phi}_{i,\text{eq}}$$

Properties:

- Single eigenfrequency
- Rhs of MDs and SMDs are different
- Neglection of mass \rightarrow SMDs
- SMDs only obtained when neglecting mass
- MDs are (in general) NOT symmetric
- Singular linear system of equations

Volterra series-based modal derivatives

$$\left(\mathbf{K}^{(1)} - (\omega_i + \omega_j)^2 \mathbf{M} \right) \tilde{\boldsymbol{\theta}}_{ij} = - \frac{1}{2} \left. \frac{\partial \mathbf{K}^{(1)}(\mathbf{q})}{\partial \eta_j(t)} \right|_{\mathbf{q}_{\text{eq}}} \boldsymbol{\phi}_i,$$

$$\left(\mathbf{K}^{(1)} - (\omega_i - \omega_j)^2 \mathbf{M} \right) \tilde{\tilde{\boldsymbol{\theta}}}_{ij} = - \frac{1}{2} \left. \frac{\partial \mathbf{K}^{(1)}(\mathbf{q})}{\partial \eta_j(t)} \right|_{\mathbf{q}_{\text{eq}}} \boldsymbol{\phi}_i,$$

Properties:

- Sum/subtraction of eigenfrequencies
- Rhs for new MDs is the same as for SMDs!
- Neglection of mass $\rightarrow \tilde{\boldsymbol{\theta}}_{ij} = \tilde{\tilde{\boldsymbol{\theta}}}_{ij} = \boldsymbol{\theta}_{s,ij}$
- Cancellation of eigenfrequencies: $\tilde{\tilde{\boldsymbol{\theta}}}_{ii} = \boldsymbol{\theta}_{s,ij}$
- New MDs are symmetric!
- Regular linear system of equations
 - only singular, if sum/subtraction of eigenfrequencies is again eigenfrequency

1.) Assess approximation quality of Volterra approximation

Compare analytical solution given by the first and second subsystem

$$\mathbf{q}_{1..2}(t) = \underbrace{\sum_{i=1}^n c_i \phi_i \cos(\omega_i t)}_{\mathbf{q}_1(t)} + \underbrace{\sum_{i=1}^n \sum_{j=1}^n \frac{1}{2} c_i c_j \left(\tilde{\theta}_{ij} \cos((\omega_i + \omega_j)t) + \tilde{\tilde{\theta}}_{ij} \cos((\omega_i - \omega_j)t) \right)}_{\mathbf{q}_2(t)}$$

with simulated solution $\mathbf{q}_{\text{sim}}(t)$ or solution given by superposing nonlinear normal modes

Reformulated analytical solution using trigonometric identities

$$\begin{aligned} \mathbf{q}_2(t) &= \sum_{i=1}^n \sum_{j=1}^n \frac{1}{2} c_i c_j \left(\tilde{\theta}_{ij} \cos((\omega_i + \omega_j)t) + \tilde{\tilde{\theta}}_{ij} \cos((\omega_i - \omega_j)t) \right) \\ &= \sum_{i=1}^n \sum_{j=1}^n c_i c_j \left(\bar{\theta}_{ij} \cos(\omega_i t) \cos(\omega_j t) - \hat{\theta}_{ij} \sin(\omega_i t) \sin(\omega_j t) \right) \end{aligned}$$

$$\bar{\theta}_{ij} = \frac{1}{2} (\tilde{\theta}_{ij} + \tilde{\tilde{\theta}}_{ij})$$

$$\hat{\theta}_{ij} = \frac{1}{2} (\tilde{\theta}_{ij} - \tilde{\tilde{\theta}}_{ij})$$



$$\mathbf{q}_{1..2}(t) = \underbrace{\sum_{i=1}^n c_i \phi_i \cos(\omega_i t)}_{\mathbf{q}_1(t)} + \underbrace{\sum_{i=1}^n \sum_{j=1}^n c_i c_j \left(\bar{\theta}_{ij} \cos(\omega_i t) \cos(\omega_j t) - \hat{\theta}_{ij} \sin(\omega_i t) \sin(\omega_j t) \right)}_{\mathbf{q}_2(t)}$$

2.) Novel quadratic manifold approaches for model order reduction

- Common quadratic manifold approach

$$\begin{aligned} \mathbf{q}(t) &\approx \sum_{i=1}^r \phi_i q_{r,i}(t) + \sum_{i=1}^r \sum_{j=1}^r \boldsymbol{\theta}_{ij} q_{r,i}(t) q_{r,j}(t) \\ &= \boldsymbol{\Phi}_r \mathbf{q}_r(t) + \boldsymbol{\Theta}_{r^2} (\mathbf{q}_r(t) \otimes \mathbf{q}_r(t)) \end{aligned}$$

Using symmetrized MDs $\boldsymbol{\theta}_{ij}$
or static MDs (SMDs) $\boldsymbol{\theta}_{s,ij}$

- Novel quadratic manifold approach (1)

$$\begin{aligned} \mathbf{q}(t) &\approx \sum_{i=1}^r \phi_i q_{r,i}(t) + \sum_{i=1}^r \sum_{j=1}^r \bar{\boldsymbol{\theta}}_{ij} q_{r,i}(t) q_{r,j}(t) \\ &= \boldsymbol{\Phi}_r \mathbf{q}_r(t) + \bar{\boldsymbol{\Theta}}_{r^2} (\mathbf{q}_r(t) \otimes \mathbf{q}_r(t)) \end{aligned}$$

Reduced coordinates:

$$\begin{aligned} q_{r,i}(t) &= c_i \cos(\omega_i t), \\ q_{r,i}(t) q_{r,j}(t) &= c_i c_j \cos(\omega_i t) \cos(\omega_j t) \end{aligned}$$

- Novel quadratic manifold approach (2)

$$\begin{aligned} \mathbf{q}(t) &\approx \sum_{i=1}^r \phi_i q_{r,i}(t) + \sum_{i=1}^r \sum_{j=1}^r \bar{\boldsymbol{\theta}}_{ij} q_{r,i}(t) q_{r,j}(t) - \frac{1}{\omega_i \omega_j} \hat{\boldsymbol{\theta}}_{ij} \dot{q}_{r,i}(t) \dot{q}_{r,j}(t) \\ &= \boldsymbol{\Phi}_r \mathbf{q}_r(t) + \bar{\boldsymbol{\Theta}}_{r^2} (\mathbf{q}_r(t) \otimes \mathbf{q}_r(t)) - \hat{\boldsymbol{\Theta}}_{r^2} (\dot{\mathbf{q}}_r(t) \otimes \dot{\mathbf{q}}_r(t)) \end{aligned}$$

$$\bar{\boldsymbol{\theta}}_{ij} = \frac{1}{2} (\tilde{\boldsymbol{\theta}}_{ij} + \tilde{\tilde{\boldsymbol{\theta}}}_{ij})$$

$$\hat{\boldsymbol{\theta}}_{ij} = \frac{1}{2} (\tilde{\boldsymbol{\theta}}_{ij} - \tilde{\tilde{\boldsymbol{\theta}}}_{ij})$$

Reduced velocities:

$$\begin{aligned} \dot{q}_{r,i}(t) &= -c_i \omega_i \sin(\omega_i t), \\ \dot{q}_{r,i}(t) \dot{q}_{r,j}(t) &= c_i c_j \omega_i \omega_j \sin(\omega_i t) \sin(\omega_j t) \end{aligned}$$

Dimensional reduction – Roadmap and Workflow

Basis augmentation

$$q(t) \approx V_{\text{aug}} q_{r,\text{aug}}(t)$$

$$V_{\text{aug}} = [\Phi_r, \Theta_{r^2}], \quad V_{\text{aug}} = [\Phi_r, \Theta_s]$$

$$V_{\text{aug}} = [\Phi_r, \tilde{\Theta}, \tilde{\tilde{\Theta}}], \quad V_{\text{aug}} = [\Phi_r, \bar{\Theta}, \hat{\Theta}]$$

Quadratic manifold

$$q(t) \approx \Phi_r q_r(t) + \Theta_{r^2} (q_r(t) \otimes q_r(t))$$

$$q(t) \approx \Phi_r q_r(t) + \bar{\Theta}_{r^2} (q_r(t) \otimes q_r(t))$$

$$q(t) \approx \Phi_r q_r(t) + \bar{\Theta}_{r^2} (q_r(t) \otimes q_r(t)) - \hat{\Theta}_{r^2} (\dot{q}_r(t) \otimes \dot{q}_r(t))$$

Application to nonlinear system

$$M\ddot{q} + f(q) = BF$$

$$y = Cq$$

Application to polynomial system

$$M\ddot{q} + K^{(1)}q + K^{(2)}(q \otimes q) + K^{(3)}(q \otimes q \otimes q) = BF$$
$$y = Cq$$

Evaluation in time-domain

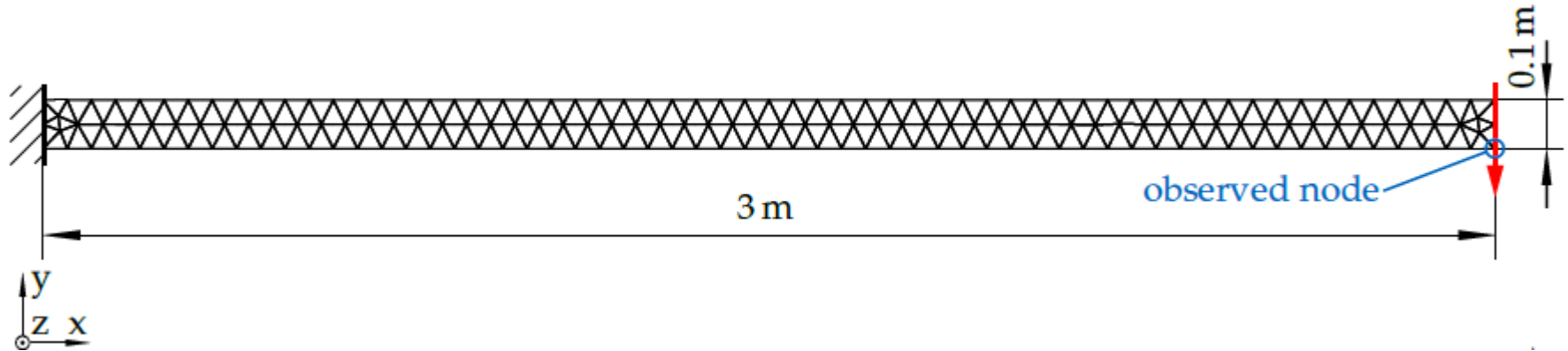
- Compare FOM and ROM via simulation runs for different inputs
- Compare ROMs with POD-ROM

Evaluation in frequency-domain

- Compute NNMs and NLFRF via shooting and path continuation
- Compare NNMs and NLFRFs of FOM and differently obtained ROMs

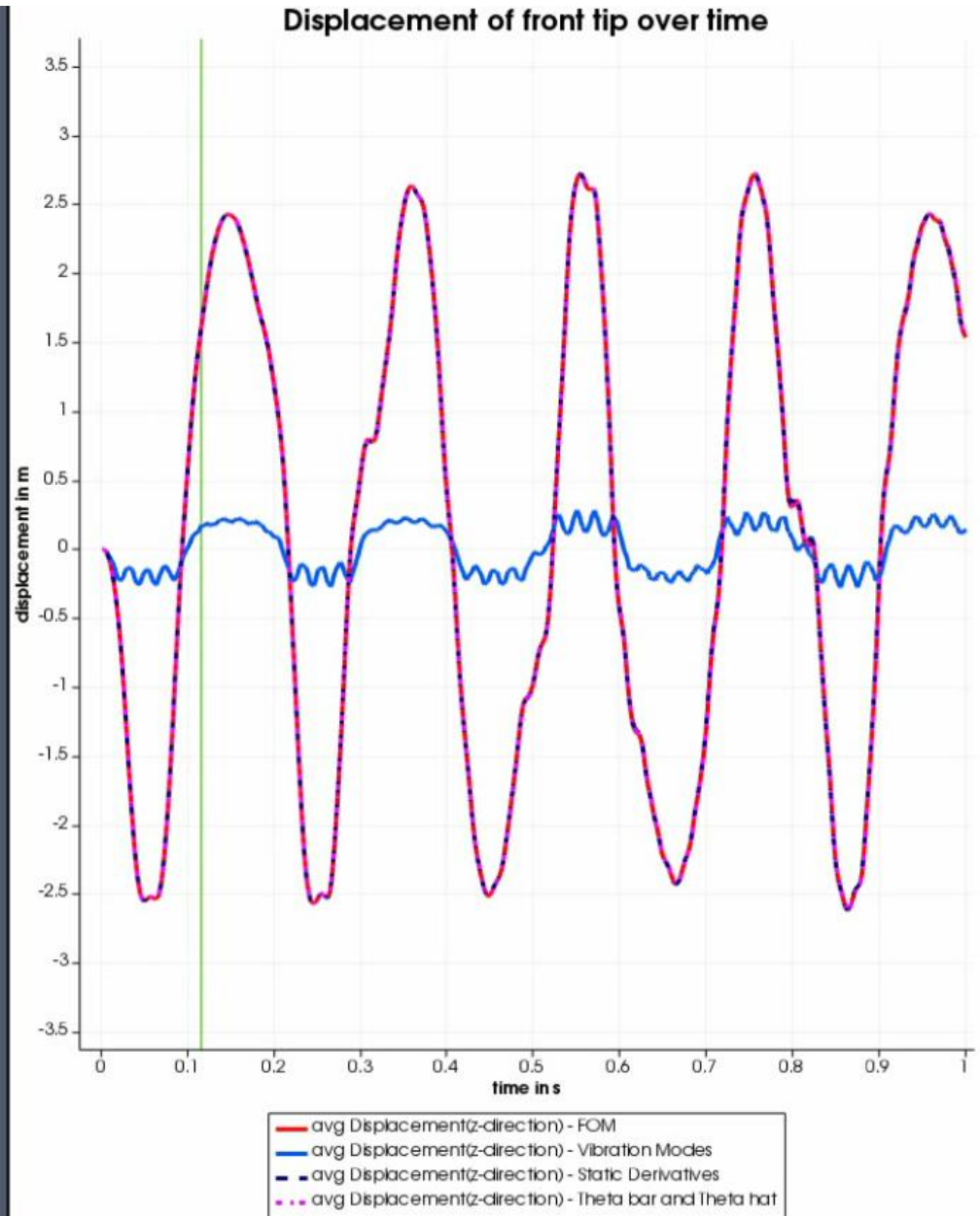
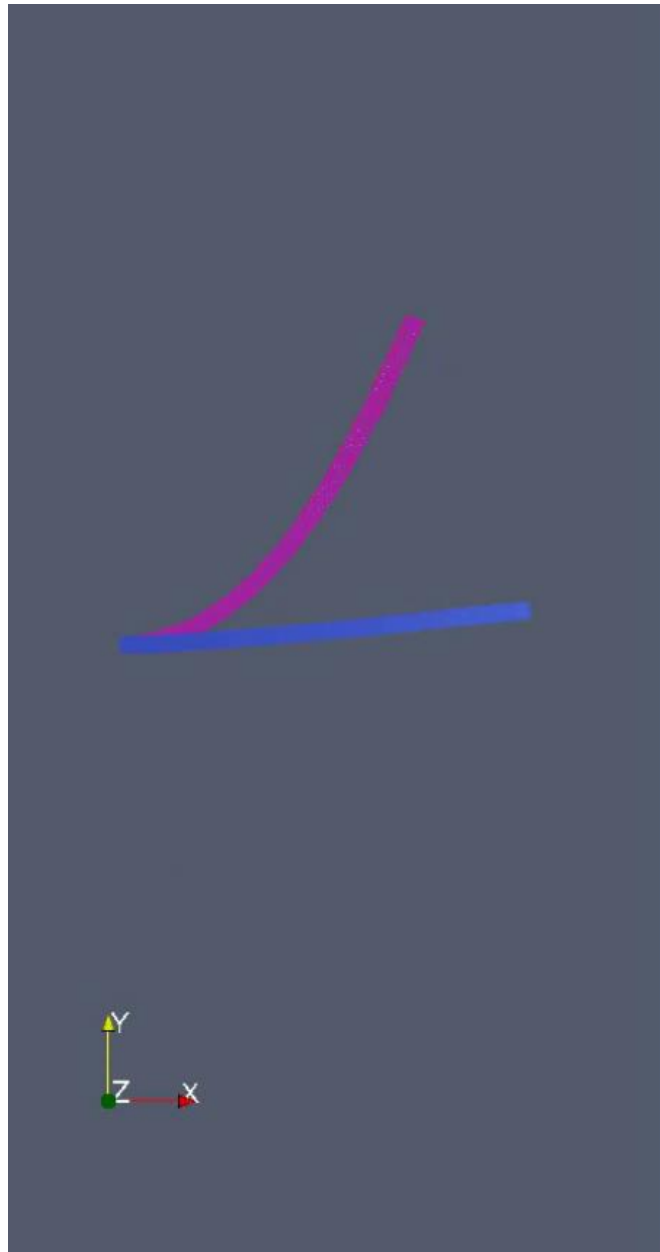
Preliminary simulation – Cantilever Beam

2D model of a cantilever beam



- 246 triangular Tri6 elements; 1224 dofs
- linear St. Venant-Kirchhoff material
- geometric nonlinear behaviour
- loading force at the tip in negative y-direction
- simulation conducted with open-source AMfe-code
- reduction via basis augmentation with new MDs

Preliminary simulation – Cantilever Beam



Summary & Outlook

Take-Home Messages:

- Model reduction with modal derivatives to capture **(geometric) nonlinear behaviour**
- Classical derivation of MDs is based on **perturbation** of the **linearized eigenvalue problem**
- Novel derivation based on **Volterra series** yields slightly different expressions for MDs
- Novel MDs are **inherently symmetric**; static derivatives can be retrieved from the new MDs
- Possible promising **applications** in nonlinear model order reduction

Ongoing / Future Work:

- Implementation of novel quadratic manifold approaches
- Validation of ROMs in time-, but also in frequency-domain (NLFRFs)
- Hyper-Reduction

Thank you for your attention!